## Some Investigations into the Dispersion Characteristics of Localized Bessel Beams

John B. Lancaster\*1, Deb Chatterjee<sup>2</sup>, and Anthony Caruso<sup>1</sup>

- Department of Physics and Astronomy, University of Missouri -Kansas City, Kansas City, MO 64110
- <sup>2</sup> School of Computing and Engineering, University of Missouri -Kansas City, Kansas City, MO 64110

An investigation into possible increases of maximal propagation distance of Bessel Beam will be presented. The Bessel Beams presented in this work will be produced by a leaky radial waveguide consisting of a capacitive sheet over a ground plane to create a normal electric field given by a truncated Bessel function (M. Ettorre and A. Grbic, *IEEE Trans. Antennas Propag.*, vol. 60, pp. 3605-3613, Aug. 2012). The required properties of leaky waveguides which best support non-diffracting waves are determined by the *n*th-order Hankel function of the first and second order, assuming TM modes. A significant feature of the localized nature of Bessel beams is the *dispersion* distance

$$Z_{max} = a\sqrt{\left(\frac{k_2}{\beta_\rho}\right)^2 - 1}. (1)$$

Beyond  $Z_{max}$ , the localized nature of the Bessel beams start to degrade. A primary objective is to seek methods by which the dispersion distance can be enhanced. In (1) the  $\beta_{\rho}$  is the phase factor of the radial propagation constant  $k_{\rho} = \beta_{\rho} - \jmath \alpha_{\rho}$ , and, a is the radius of the circular aperture. The aperture radiates into medium # 2 defined by its intrinsic wavenumber  $k_2$ . Obviously, it follows from (1), that if  $\beta_{\rho}$  is decreased, and the circular aperture radius a is increased, then the dispersion distance can be enhanced.

To that end, by invoking the conditions of producing the pure Bessel  $(J_o(z))$  beam in the circular aperture, use is made of the following Hankel large argument forms, that read:

$$H_{\circ}^{(1),(2)}(z) \approx \sqrt{\frac{2}{\pi z}} [\mathcal{P}(0,z) \pm j\mathcal{Q}(0,z)] e^{\pm j(z - \frac{\pi}{4})}.$$
 (2)

In (2),  $\mathcal{P}(0,z) \approx 1$ , and  $\mathcal{Q}(0,z) \approx \frac{-1}{8z}$ . The argument  $z = k_{\rho}\rho$  in (2). Invoking the conditions of producing the pure Bessel beam one then obtains the dispersion relationship

$$\tan(z - \frac{\pi}{4}) = 8z. \tag{3}$$

It is thus obvious that solutions to (3) shall lead to real values of  $\beta_{\rho}$ . This approach is different from the earlier work of Ettore and Grbic, op. cit. There could be several forms of the dispersion relationships that are distinct from (3), using different approximations to the Hankel functions, and hence a host of solutions for the radial propagation constant are feasible. These aspects shall be discussed in the view of the results from full-wave EM simulations for characterizing the dispersion distance of localized Bessel beams.