

Integral equation methods for solving problems of scattering on an unclosed cone structure

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The method based on using the integral Kontorovich-Lebedev transforms and the semi-inversion method is a successful one for studying boundary electrodynamics problems with single periodical slotted cone geometry (there is one cone strip in the period). Taking into consideration a multi-element cone structure (there are several strips in the period) makes applying the above mentioned method extremely tedious. The task of this work is to find a method that uses singular integral equation to solve an excitation problem for a periodic cone structure with any number of strips (S) in the period.

Let a semi-infinite perfectly conducting infinitely thin circular cone Σ with N equally spaced slots cut along the generatrices be excited by a radial harmonic electric or magnetic dipole. The angular width of the slots and the period of the cone are d and $l = 2p/N$ respectively. By virtue of introducing Debye potentials and the integral Kontorovich-Lebedev transforms the electrodynamics boundary problem is reduced to solving dual series equations for unknown Fourier coefficients $\hat{y}_{m,n}^{(j)}$ of electromagnetic field components. The dual series equations are equivalent to singular integral equations like these:

1) for electrical dipole excitation, $j=1$,

$$\frac{1}{2p} \int_S \hat{F}(\mathbf{y} - \mathbf{a}) \Phi_1(\mathbf{a}) d\mathbf{a} + \frac{1}{2p} \int \left[\hat{A}_{nt}^{(1)} - \hat{K}_1(\mathbf{y} - \mathbf{a}) \right] \Phi_1(\mathbf{a}) d\mathbf{a} = e^{im_0 \mathbf{y}},$$

$$\mathbf{y} \in S : \frac{pd}{l} < |\mathbf{y}| \leq p, \hat{F}(\mathbf{y} - \mathbf{a}) = \sum_{n \neq 0} \frac{1}{N(n+\mathbf{n})} \frac{|n|}{n} e^{in(\mathbf{y}-\mathbf{a})},$$

$$\hat{K}_1(\mathbf{y} - \mathbf{a}) = \sum_{n \neq 0} \frac{1}{N(n+\mathbf{n})} \frac{|n|}{n} \hat{\mathbf{e}}_{m,n}^{(1)} e^{in(\mathbf{y}-\mathbf{a})};$$

2) for magnetic dipole excitation, $j=2$,

$$\frac{1}{p} \int_S \frac{\Phi_2(\mathbf{x})}{\mathbf{x} - \mathbf{y}_1} d\mathbf{x} + \frac{1}{p} \int K(\mathbf{x} - \mathbf{y}_1) \Phi_2(\mathbf{x}) d\mathbf{x} = ie^{im_0 \mathbf{y}_1}, \mathbf{y}_1 \in S,$$

$$K(\mathbf{x} - \mathbf{y}_1) = \frac{1}{2} \text{ctg} \frac{\mathbf{x} - \mathbf{y}_1}{2} - \frac{1}{\mathbf{x} - \mathbf{y}_1} - \frac{i}{2N} \left(\frac{pe^{inx}}{\sin pn} - \frac{1}{n} \right) \frac{1}{A_{nt}^{(2)}} - \frac{i}{2} \sum_{n \neq 0} \frac{|n|}{n} \hat{\mathbf{e}}_{m,n}^{(2)} e^{in(\mathbf{y}_1 - \mathbf{x})};$$

$$\Phi_j(\mathbf{y}) = \sum_{n=-\infty}^{+\infty} \hat{y}_{m,n}^{(j)} e^{iny}, \mathbf{y} \in [-p, p], \hat{\mathbf{e}}_{m,n}^{(j)} = O\left(\frac{\sin^2 \mathbf{g}}{N^2(n+\mathbf{n})^2}\right), -1/2 \leq \mathbf{n} < 1/2,$$

$m, m_0 \in Z$, $\hat{A}_{nt}^{(1)}, A_{nt}^{(2)}$ are known.

These integral equations can be solved numerically by the discrete singularities method.